

# Weakness of gravity as illusion which hides true path to unification of gravity with particle physics

Alexander Burinskii

NSI, Russian Academy of Sciences,

B. Tuskaya 52 Moscow 115191 Russia, email: bur@ibrae.ac.ru

Essay written for the Gravity Research Foundation 2017 Awards for

Essays on Gravitation

March 30, 2017

## **Abstract**

Well known weakness of Gravity in particle physics is an illusion caused by underestimation of the role of spin in gravity. Relativistic rotation is inseparable from spin, which for elementary particles is extremely high and exceeds mass on 20-22 orders (in units  $c = G = m = \hbar = 1$ ). Such a huge spin generates frame-dragging that distorts space much stronger than mass, and effective scale of gravitational interaction is shifted from Planck to Compton distances. We show that compatibility between gravity and quantum theory can be achieved without modifications of Einstein-Maxwell equations, by coupling to a supersymmetric Higgs model of symmetry breaking and forming a nonperturbative super-bag solution, which generates a gravity-free Compton zone necessary for consistent work of quantum theory. Super-bag is naturally upgraded to Wess-Zumino supersymmetric QED model, forming a bridge to perturbative formalism of conventional QED.

As is known, Quantum theory and Gravity cannot be combined in a unified theory. Gravity refuses pointlike, structureless quantum particles, requiring extended field structure for right side of Einstein equations,  $G_{\mu\nu} = 8\pi T_{\mu\nu}$ .

Revolutionary step for unification was made in superstring theory – transition to extended stringlike objects, however, “...*Since 1974 superstring theory stopped to be considered as particle physics...* ” and “... *a realistic model of elementary particles still appears to be a distant dream ...* ”, J. Schwarz [1],

One of reason was the choice of Planck scale as universal scale for all unifications, including gravity. Attempt to bring gravitational scale close to the weak scale was considered, in the braneworld scenario, where the weakness of the localized 4d gravity was explained by its “leaks” into higher-dimensional bulk. Braneworld mechanism allowed to realize ideas of the superstring theory for any numbers of extra dimensions [2].

Alternative approach was related with solitons – nonperturbative 4D solutions of the nonlinear field models, in particular solutions to low energy string theory [3, 4, 5]. This approach, being akin to Higgs mechanism of symmetry breaking, is matched with nonperturbative approach to electroweak sector of Standard Model, where the most known are the Nielsen-Olesen model of dual string based on the Landau-Ginzburg (LG) field model for superconducting media, and the famous MIT and SLAC bag models [6, 7, 8] which are similar to solitons, but being soft, deformable and oscillating, acquire many properties of the string models. The question on consistency with gravity is not discussed usually for the solitonic models, as it is conventionally assumed that gravity is weak and not essential at scale of electroweak interactions (see, for example, [9, 10])

We claim that assumption on weakness of gravity is an illusion, related with underestimation of the role of spin in gravity. In relativistic theory spin is inseparable from rotation, and created by spin invariant effect of gravitational frame-dragging [11] (supported by Probe B experiment), or Lense-Thirring effect in Kerr geometry [12], distorts space along with mass.

Spin of elementary particles is extremely high. In particular, for electron spin/mass ratio is about  $10^{22}$  (in dimensionless units  $G = c = \hbar = 1$ ), and its influence becomes so strong that conflict with quantum theory is shifted from Planck to Compton scale. Similar to Cosmology where giant masses turn gravity into a main force, GIANT SPIN of particles MAKES GRAVITY STRONG and crucial by its interplay with quantum theory.

The spinning Kerr-Newman (KN) solution [12, 13] is of particular interest, since it has gyromagnetic ratio  $g = 2$ , corresponding to Dirac theory of electron [14], and structure of the KN solution with such a huge spin sheds light on the reason of conflict between gravity and quantum theory and points out the way for its resolution.

Metric of KN solution in the Kerr-Schild form is [13]

$$g_{\mu\nu} = \eta_{\mu\nu} + 2Hk_{\mu}k_{\nu}, \quad (1)$$

where  $\eta_{\mu\nu}$  is metric of auxiliary Minkowski space  $M^4$ , (signature  $(-+++)$ ), where scalar function

$$H_{KN} = \frac{mr - e^2/2}{r^2 + a^2 \cos^2 \theta}, \quad (2)$$

is given in oblate spheroidal coordinates  $r$  and  $\theta$ , determined by transformations [13],

$$x + iy = (r + ia) \exp\{i\phi_K\} \sin \theta, \quad z = r \cos \theta, \quad \rho = r - t. \quad (3)$$

The null field  $k_{\mu}(x)$ , ( $k_{\mu}k^{\mu} = 0$ ) determines directions of dragging of space, Fig.1. Lense-Thirring effect creates vortex of congruence  $k_{\mu}(x)$ , which for ultra-high spin,  $a = J/m \gg m$ , becomes so strong that BH horizons disappear, and  $k^{\mu}(x)$  focus on singular ring  $r = 0, \cos \theta = 0$ , forming branch line of space into two sheets,  $g_{\mu\nu}^{\pm} = \eta_{\mu\nu} + 2Hk_{\mu}^{\pm}k_{\nu}^{\pm}$ , defined by ingoing  $k_{\nu}^{-}$  and outgoing congruence  $k_{\nu}^{+}$ .

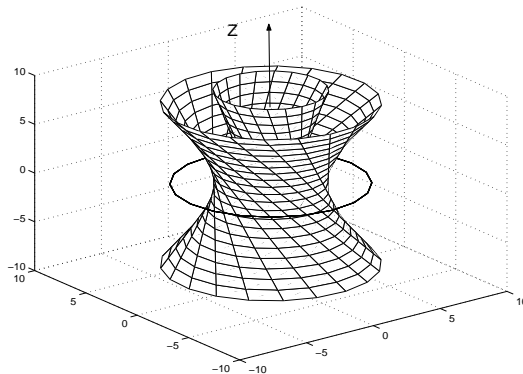


Figure 1: The lightlike Kerr congruence  $k^{\mu}$  determines space-dragging caused by mass and spin. Two sheets of Kerr metric correspond to  $r < 0$ , and  $r > 0$ .

Quantum theory requires flat space, at least in Compton zone, but electron spin  $J = \hbar/2$  exceeds mass  $m$  about 22 orders, which breaks space topologically, creating singular ring of Compton radius  $a = \hbar/2m$ . Singularity is signal to new physics. Usually, it is considered as signal to modify gravity. We suggest alternative solution based on *supersymmetry*, which expels gravitational field from Compton zone of spinning particle, similar to expulsion of electromagnetic field from superconductor. Supersymmetric bag model, [15, 16, 17], realizes such expulsion of gravity and electromagnetic field, forming tree zones:

- (**I**) – flat quantum interior,
- (**E**) – external zone with exact KN solution,
- (**R**) – zone of transition from (**I**) to (**E**).

For the giant values of spin, these demands become so restrictive that structure of bag is determined almost unambiguously.

Surface ( $R$ ) is defined by the continuous transition of KN solution to Minkowski interior of the bag, (C. López [18]). According (1) and (2), zone (**R**) corresponds to

$$H_{KN}(r) = 0. \quad (4)$$

which gives

$$r = R = \frac{e^2}{2m}, \quad (5)$$

and relations (4) and (3) show that bag takes form of a disk with thickness  $R$  and radius  $r_c = \sqrt{R^2 + a^2}$ , Fig.2.

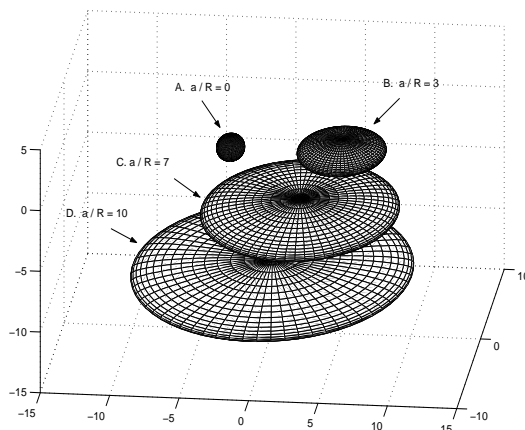


Figure 2: Shape of disk for different  $a = J/m$ : (A) -  $a/R = 0$ , (B)-  $a/R = 3$ ; (C) -  $a/R = 7$ ; and (D) -  $a/R = 10$ .

To satisfy (**I**),(**E**),(**R**), it is natural to use Higgs mechanism of symmetry breaking which is used in many nonperturbative electroweak models, [19], and also in the MIT and SLAC bag models [6, 7]. The corresponding Lagrangian is also known as Landau-Ginzburg (LG) field model for superconducting phase transitions. The famous Nielsen-Olesen (NO) model for vortex string in superconducting media, [20], is based on LG Lagrangian

$$\mathcal{L}_{NO} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{2}(\mathcal{D}_\mu\Phi)(\mathcal{D}^\mu\Phi)^* - V(|\Phi|), \quad (6)$$

where  $\mathcal{D}_\mu = \nabla_\mu + ieA_\mu$  are  $U(1)$  covariant derivatives,  $F_{\mu\nu} = A_{\mu,\nu} - A_{\nu,\mu}$  the corresponding field strength, and potential  $V$  has typical form

$$V = \lambda(\Phi^\dagger\Phi - \eta^2)^2, \quad (7)$$

where  $\eta$  is condensate of Higgs field  $\Phi$ ,  $\eta = \langle |\Phi| \rangle$ .

However, potential (7) distorts external KN solution, placing Higgs field in zone (E). It turns out that conditions **(I)**, **(E)**, **(R)** are satisfied by supersymmetric LG model with three Higgs-like fields, [21]  $(H, Z, \Sigma) \equiv (\Phi_1, \Phi_2, \Phi_3)$ .

Corresponding Lagrangian differs from (6) only by summation over the fields  $\Phi_i, i = 1, 2, 3$ , while the potential  $V$  is changed very essentially, and formed from a superpotential function  $W(\Phi_i)$ , [22]

$$V(r) = \sum_i F_i F_i^*, \quad F_i = \partial W / \partial \Phi_i \equiv \partial_i W, \quad (8)$$

where

$$W(\Phi_i, \bar{\Phi}_i) = Z(\Sigma\bar{\Sigma} - \eta^2) + (Z + \mu)H\bar{H}, \quad (9)$$

The conditions  $F_i = \partial_i W = 0$  determine *two vacuum states*,  $V = 0$  :

**(I)** - internal:  $r < R - \delta$ , where Higgs field  $|H| = \eta$ ,

**(E)** - external:  $r > R + \delta$ , where Higgs field  $H = 0$ ,

separated by *zone of phase transition* **(R)**,  $V > 0$ , in correspondence with **(I)**, **(E)**, **(R)**.

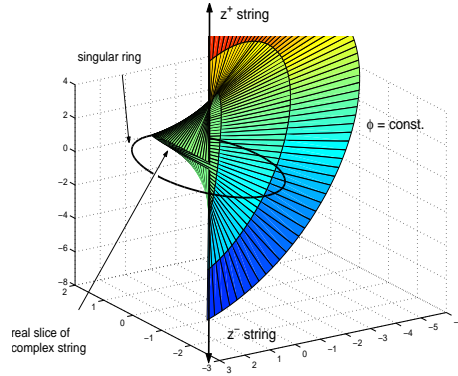


Figure 3: Rotational dragging of potential  $A^\mu$ , forms Wilson loop along border of disk.

Bag models with potential (7) form "cavity in superconductor", [8]. Supersymmetric potential (8)-(9) concentrates Higgs field in zone **(I)**, forming supersymmetric and superconducting vacuum state inside bag, where (6) gives equation

$$\square A_\mu = J_\mu = e|H|^2(\chi_{,\mu} + eA_\mu), \quad (10)$$

showing that current  $J_\mu = 0$ , is concentrated (as usual, [23]) in surface layer of superconducting disk.

Near boundary of disk  $r = R = e^2/2m$ ,  $\cos\theta = 0$ , vector-potential  $A_\mu$  is dragged by Kerr congruence (Fig.2), forming closed Wilson loop along singular ring. *It has remarkable consequence – quantization of angular momentum*, [16, 17, 24],  $J = n/2$ ,  $n = 1, 2, 3, \dots$

Bag models take intermediate position between strings and solitons [25, 26, 27]. Similar to solitons, they are nonperturbative solutions of the Higgs field model, but they have several specific features, in particular, flexibility and ability to create string-like structures. Under rotation, bags are deformed and take shape of stringy flux-tube joining the quark-antiquark pair [6, 26], or toroidal string [7, 27, 28].

Spinning gravitational field sets shape of bag according **(R)**, and circular string is formed on the boundary of the disk, closely to Kerr singular ring (Fig.4A). So, the string is really formed by singular ring and regularized by bag boundary. The assumption, that Kerr singular ring is similar to NO model of dual string was done long ago in [29, 30], where it was noted that excitations of the KN solution create traveling waves along the Kerr ring. Later, it was obtained in [5] close connection of the Kerr singular ring with the Sen fundamental string solution to low energy string theory [31], and other relations of Kerr geometry with string theory [32]. String admits traveling waves, which deform position of the bag boundary according **(R)**, creating a circulating lightlike node, where surface of the deformed bag touches the Kerr singular ring, creating a circulating lightlike singular pole, which can be associated with a confined quark, Fig.4B, and super-bag turns into a single "bag-string-quark" system, analog of  $D2 - D1 - D0$ -brane of string-Mtheory.

In turn, capture of quarks, is one more special feature of the bag model, requiring consistent implementation of the Dirac equation [6, 7, 27, 28]. In KN geometry, it is defined according to famous Kerr Theorem [13, 33] which defines the shear-free Kerr congruences in twistor terms, and gives two roots  $Y^\pm$  for projective spinor coordinate

$$Y = \phi_1/\phi_0, \quad (11)$$

which is equivalent to two-component Weyl spinor  $\phi_\alpha$ , and defines the null direction as

$$k_\mu = \bar{\phi}_{\dot{\alpha}} \sigma_\mu^{\dot{\alpha}\alpha} \phi_\alpha. \quad (12)$$

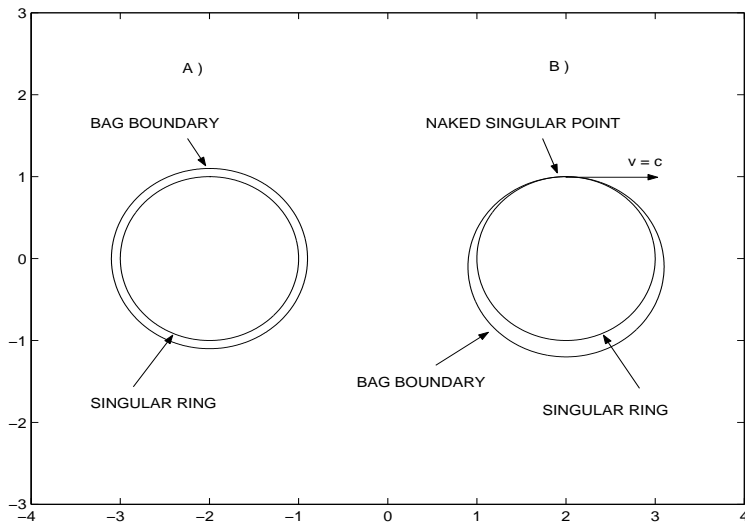


Figure 4: Regularization of the KN string. Boundary of bag fixes cut-off  $R = r_e$  for the Kerr singular ring: A) exact KN solution, B) KN solution excited by the lowest traveling mode creates singular pole.

As it was shown in [34, 15], two roots of the Kerr theorem  $Y^\pm$  give us two Weyl spinors,  $\phi_\alpha$  and  $\bar{\chi}^{\dot{\alpha}}$ , consistent with ingoing and outgoing KN solutions in zone **(E)**. Inside the bag, these solutions unite, forming the Dirac spinor

$$\Psi = \begin{pmatrix} \phi_\alpha \\ \bar{\chi}^{\dot{\alpha}} \end{pmatrix} \quad (13)$$

which gets mass through Yukawa coupling to condensate of the Higgs field. Here we meet third feature of the bag model – emergence of the position-dependent mass term  $m = G|\Phi|$ , which is determined by spacetime distribution of the Higgs field.

The Dirac equation,  $\gamma^\mu \partial_\mu \Psi = m\Psi$ , which is massive inside the bag, turns out to be massless, and splits into two independent massless equations

$$\sigma_{\alpha\dot{\alpha}}^\mu i\partial_\mu \bar{\chi}^{\dot{\alpha}} = 0, \quad \bar{\sigma}^{\mu\dot{\alpha}\alpha} i\partial_\mu \phi_\alpha = 0, \quad (14)$$

outside the bag, corresponding to the left-handed and right-handed “electron-type leptons” of the Glashow-Salam-Weinberg model [35].

Finally, Super-Bag can be naturally upgraded to Wess-Zumino SuperQED model, [36], revealing connections between the non-perturbative solutions of the supersymmetric LG model and the conventional perturbative technics used in QED.

**Conclusion:**

Weakness of Gravity is delusion caused by the underestimation of huge impact of spin on space-time metric and topology. Disposal of this delusion opens a supersymmetric way to unify Gravity with particle physics.



## References

- [1] J. Schwarz, *The Early History of String Theory and Supersymmetry*, CALT-68-2858 [arXiv:1201.0981]
- [2] N. Arkani-Hamed<sup>1</sup>, S. Dimopoulos, G. Dvali<sup>4</sup> and N. Kaloper, Infinitely Large New Dimensions, *Phys.Rev.Lett.* **84** 586 (2000) DOI: 10.1103/PhysRevLett.84.586 [hep-th/9907209].
- [3] A. Dabholkar, J. P. Gauntlett, J. A. Harvey, D. Waldram, Strings as Solitons & Black Holes as Strings, *Nucl.Phys.* **B474**, 85 (1996) [arXiv:hep-th/9511053].
- [4] A. Sen, Macroscopic Charged Heterotic String. *Nucl.Phys.* **B 388** 457 (1992), [arXiv:hep-th/9206016].
- [5] Burinskii A., Some properties of the Kerr solution to low-energy string theory. *Phys. Rev. D* **52** 5826 (1995), [arXiv:hep-th/9504139].
- [6] Chodos A. et al. New extended model of hadrons. *Phys. Rev.* **D 9**, 3471 (1974).
- [7] Bardeen W. A. et al., Heavy quarks and strong binding: A field theory of hadron structure. *Phys. Rev.* **D 11**, 1094 (1974).
- [8] R. Dashen, B. Hasslacher, and A. Neveu, Nonperturbative methods and extended-hadron models in field theory: I,II, III. *Phys. Rev.* **D 10**, 4114 (1974), *ibid.* 4130 (1974), *ibid* 4138 (1974) .
- [9] J.M. Maldacena, *Int.J.Mod.Phys.* **A15S1** , 840 (2000)
- [10] J. C. Baez, Higher Dimensional Algebra and Planck Scale Physics, In *Physics Meets Philosophy at the Planck Length*, eds. Craig Callender and Nick Huggett, Cambridge U. Press, Cambridge, 2001, pp. 177-195, [arXiv:gr-qc/9902017]
- [11] Ch. W. Misner, K. S. Thorne, J.A. Wheeler, *Gravitation*, Part 3, San Francisco: W. H. Freeman, ISBN 978-0-7167-0344-0
- [12] R.P. Kerr, Gravitational field of a spinning mass as an example of algebraically special metrics. *Phys.Rev.Letters* **11**, 237 (1963).
- [13] Debney G. C., Kerr R. P. and Schild A. Solutions of the Einstein and Einstein-Maxwell equations. *J. Math. Phys.* **10** 1842 (1969).

- [14] Carter B., “Global structure of the Kerr family of gravitational fields,” *Phys. Rev.* **174**, 1559 (1968). *Phys. Rev.* **174**, 1559 (1968).
- [15] A. Burinskii, *Gravitating lepton bag model*, *JETP (Zh. Eksp. Teor. Fiz.)* **148**(8) (2015) 228, [arXiv:1505.03439].  
A. Burinskii, *Int. J. Mod. Phys. Source of the Kerr-Newman solution as a gravitating bag model: 50 years of the problem of the source of the Kerr solution*, *Int. J. Mod. Phys.* **A31**, 1641002 (2016).
- [16] A. Burinskii, *Stability of the lepton bag model based on the Kerr-Newman solution*, *JETP (Zh. Eksp. Teor. Fiz.)* **148** (2015) 937.
- [17] A. Burinskii, *Source of the Kerr-Newman solution as a supersymmetric domain-wall bubble: 50 years of the problem*, *Phys Lett. B*, 754 (2016) 99.
- [18] López C.A. , *An Extended Model Of The Electron In General Relativity*, *Phys. Rev. D* **30**, 313 (1984).
- [19] Coleman S., Q-Balls, *Nuclear Physics B* **262** (2) 263 (1985).  
Rosen G., Particlelike Solutions to Nonlinear Complex Scalar Field Theories with Positive-Definite Energy Densities. *J. of Math. Phys.* **9** (7) 996 (1968), doi:10.1063/1.1664693 .  
A. Achúcarro and T. Vachaspati, Semilocal and Electroweak Strings, *Phys.Rept.* **327** 347 (2000), [arXiv:hep-ph/9904229].  
A. Kusenko, Solitons in the supersymmetric extensions of the standard model, *Phys.Lett.* **B405** 108 (1997).  
G. Dvali, A. Kusenko and M. Shaposhnikov, “New physics in a nutshell, or Q-ball as a power plant”, *Phys.Lett.* **B417**, 99 (1998).  
Volkov M. and Wöhrner E., Spinning Q-balls, *Phys.Rev. D* **66**, 085003 (2002).
- [20] H. B. Nielsen and P. Olesen, Vortex-line models for dual strings, *Nucl. Phys.* **B 61**, 45, (1973).
- [21] J. R. Morris, *Phys. Rev. D* **53** 2078 (1996) [arXiv:hep-ph/9511293].
- [22] Wess J. and Bagger J., *Supersymmetry and Supergravity*, Princeton Univ. Press, New Jersey, 1983.

- [23] L.D. Landau and E.M. Lifshitz, *Electrodynamics of Continuous Media (Volume 8 of A Course of Theoretical Physics)*, Pergamon Press, 1960.
- [24] Burinskii A., Regularized Kerr-Newman Solution as a Gravitating Soliton. *J. Phys. A: Math. Theor.* **43**, 392001 (2010), [arXiv: 1003.2928].  
 Burinskii A., Kerr-Newman electron as spinning soliton, *Int J. of Mod. Phys. A* **29**, 1450133 (2014), [arXiv:1410.2888].
- [25] Giles R.C., Semiclassical dynamics of the "SLAC bag", *Phys. Rev.* **D 70**, 1670 (1976).
- [26] K. Johnson and C. B. Thorn, Stringlike solutions of the bag model, *Phys. Rev.* **D 13**, 1934 (1976).
- [27] S.-H. H. Tye, Quark-binding string, *Phys. Rev.* **D 13**, 3416 (1976)
- [28] R. C. Giles and S.-H. H. Tye, Quantum dynamics of a quark-binding bubble in two space and one time dimensions, *Phys. Rev.* **D 13**, 1690 (1976)
- [29] Burinskii A.Ya., Microgeons with spin. *Sov. Phys. JETP* **39** 193 (1974).
- [30] Ivanenko D.D. and Burinskii A.Ya., Gravitational strings in the models of elementary particles, *Izv. Vuz. Fiz.*, **5**, 135 (1974).
- [31] A. Sen, Rotating charged black hole solution in heterotic string theory, *Phys. Rev. Lett.* **69**, 1006, (1992).
- [32] A. Burinskii, Twistorial analyticity and three stringy systems of the Kerr spinning particle, *Phys. Rev.* **D 70**, 086006 (2004) [arXiv:hep-th/0406063].  
 A. Burinskii, Orientifold D-String in the Source of the Kerr Spinning Particle, *Phys. Rev.* **D 68** (2003) 105004 [arXiv:hep-th/0308096].  
 A. Burinskii, Stringlike structures in the real and complex Kerr-Schild geometry, *Journal of Physics: Conference Series* **532** 012004 (2014), [arXiv:1410.2462].  
 A. Burinskii, Kerr spinning particle, strings, and superparticle models, *Phys. Rev.* **D 57**, 2392 (1998).
- [33] R. Penrose, Twistor Algebra, *J. Math. Phys.* **8** 345 (1967),  
 R. Penrose and W. Rindler, *Spinors and Space-time, Vol. 2: Spinor and twistor methods in space-time geometry*, Cambridge University Press, Cambridge U.K. (1986), pg. 501.

- [34] A. Burinskii, Emergence of the Dirac Equation in the Solitonic Source of the Kerr Spinning Particle, *Grav. and Cosmol.* **21**(1) 28 (2015), [arXiv:1404.5947].
- [35] S. Weinberg, A model of leptons, *Phys.Rev.Lett.* **19**, 1264 (1967).
- [36] A. Burinskii, New Path to Unification of Gravity with Particle Physics, arXiv:1701.01025 (submitted to Phys.Rev.D).
- [37] P. A. M. Dirac, An Extensible Model of the Electron. *Proc. R. Soc. Lond. A* **268**, 57 (1962).